

The behavior of the electron cyclotron drift instability (ECDI) inside Hall Effect Thruster with a fluid code

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1) CERFACS, Toulouse FRANCE

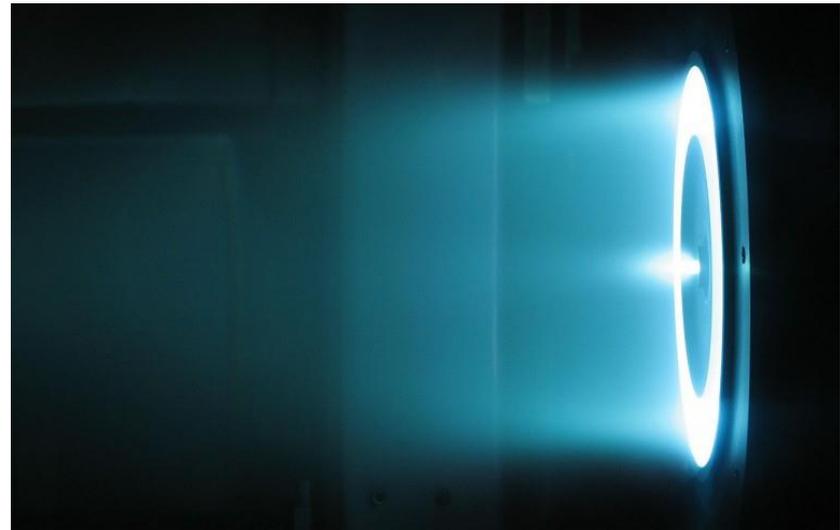
2) SAFRAN Aircraft Engines, Vernon, FRANCE

Hall effect thrusters: the industrial point of view

SAFRAN industrializes a complete range of plasma Hall effect thrusters for low power electric propulsion:



The PPS-1350 Hall thruster (SAFRAN)



The PPS-5000 Hall thruster (SAFRAN)

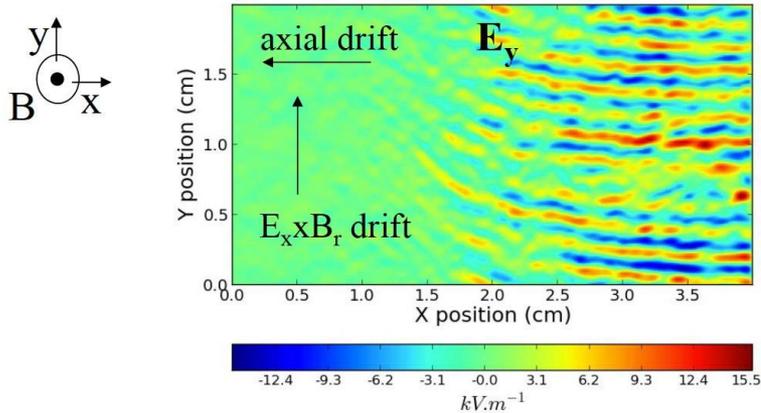
➡ Long and costly qualifications in vacuum test facilities

- ◆ Until now, no industrial numerical solver is mature enough to help them in the conception and understanding of Hall effect thrusters

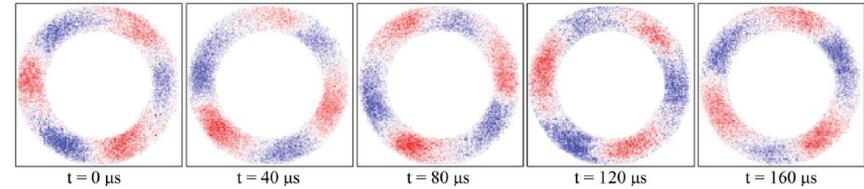
The big complexity of Hall effect thruster plasma flows

- ◆ Complex plasma phenomena are highlighted through experiments ...

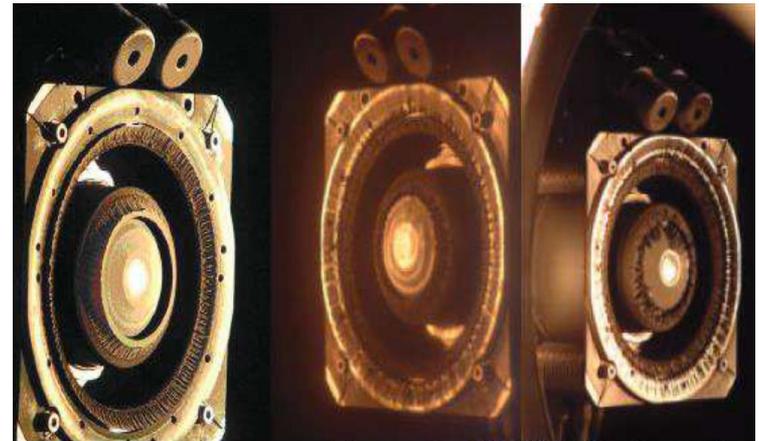
- ◆ ... or through simplified numerical problems ...



2D structure of the fluctuating azimuthal component of the electric field in a Hall thruster⁽¹⁾



Rotating spoke instabilities of plasma density on a 6 kW Hall Thruster⁽²⁾

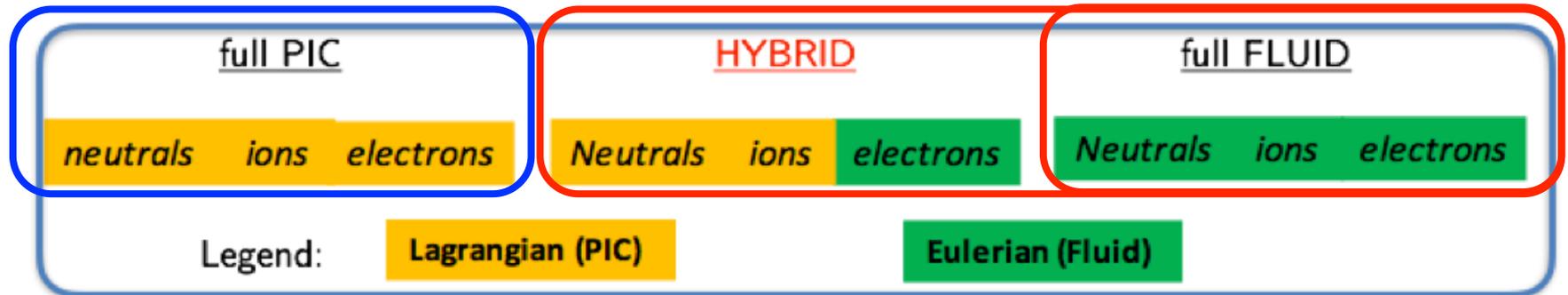


Ageing of the PPS-1350 Hall thruster (SAFRAN)

- ◆ ... but are still not fully understood. They can decrease thruster efficiency and cause an anormal erosion of ceramics

A numerical solver to model Hall-Effect thruster's flows

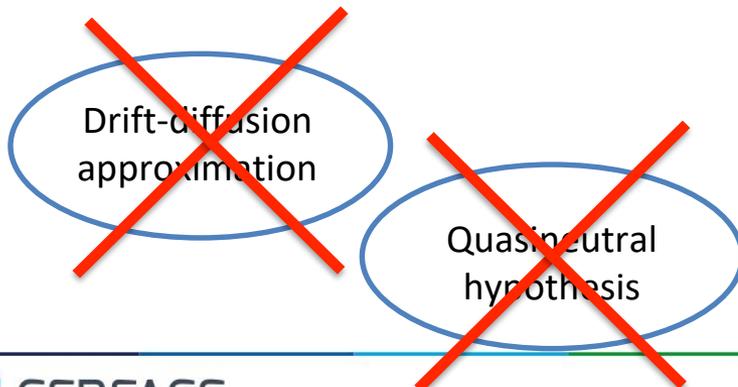
Objectives: Develop a **3D unstructured massively parallel** PIC/fluid solver to model with accuracy the plasma behavior inside a Hall Effect Thruster (HET) in collaboration with SAFRAN : **AVIP**



Too costly for an industrial complex simulation
Reference solver for fluid simulations

- More hypothesis / models
More adapted to conception simulations

In most fluid solvers:



AVIP-Fluid: Equations

- Fluid equation system for **ions and electrons** based on a 10-moment model :

$$\begin{aligned}
 \partial_t n_e + \nabla \cdot (n_e \vec{u}_e) &= n_e n_n K_{ioniz} \quad , \quad \partial_t n_i + \nabla \cdot (n_i \vec{u}_i) = n_e n_n K_{ioniz} \\
 \partial_t (m_e n_e \vec{u}_e) + \nabla \cdot (m_e n_e \vec{u}_e \vec{u}_e + k_B T_e n_e \vec{I}) &= -en_e (\vec{E} + \vec{u}_e \times \vec{B}) - K_{en} n_n n_e m_e (\vec{u}_e - \vec{u}_n) \\
 \partial_t (m_i n_i \vec{u}_i) + \nabla \cdot (m_i n_i \vec{u}_i \vec{u}_i + k_B T_i n_i \vec{I}) &= en_i (\vec{E} + \vec{u}_i \times \vec{B}) + K_{in} n_n n_i m_i (\vec{u}_i - \vec{u}_n) \\
 \partial_t (m_e n_e E_e) + \nabla \cdot \left(\left(\frac{1}{2} m_e n_e \vec{u}_e^2 + \frac{\gamma}{\gamma-1} P_e \right) \cdot \vec{u}_e + \vec{Q}_e \right) &= -en_e \vec{E} \cdot \vec{u}_e - S_{ioniz,e}^2 - S_{exc,e}^2 - S_{en,e}^2 \\
 \partial_t (m_i n_i E_i) + \nabla \cdot \left(\left(\frac{1}{2} m_i n_i \vec{u}_i^2 + \frac{\gamma}{\gamma-1} P_i \right) \cdot \vec{u}_i \right) &= en_i \vec{E} \cdot \vec{u}_i + S_{ioniz,i}^2 - S_{in,i}^2
 \end{aligned}$$

with :

- $q_\alpha n_\alpha (\vec{E} + \vec{u}_\alpha \times \vec{B})$ the Lorentz force with \vec{E} and \vec{B} the electromagnetic fields, \vec{B} is supposed constant, a Poisson equation solves

the variation of the electric potential:

$$\Delta \phi = \frac{e}{\epsilon_0} (n_i - n_e)$$

- \vec{Q}_e the electron heat flux is first neglected:

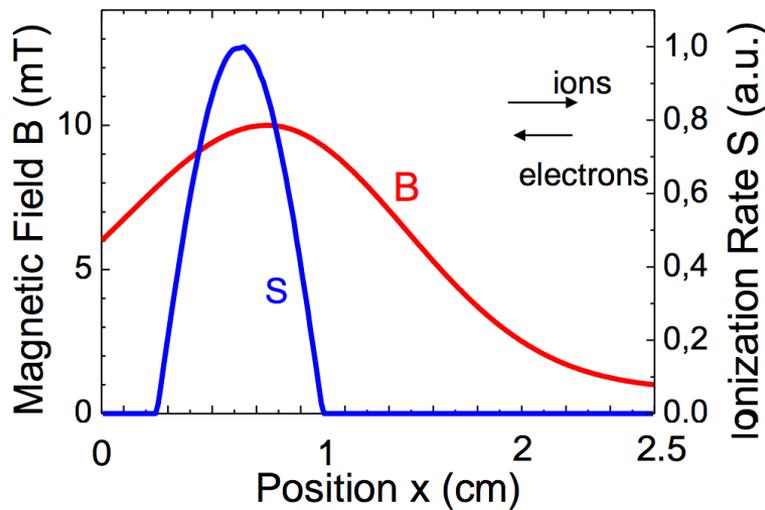
$$\vec{Q}_e = 0$$

- A mass continuity equation for neutrals at constant speed:

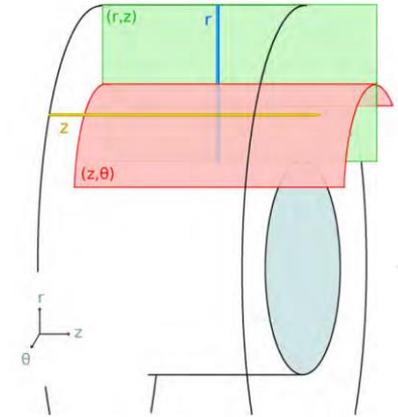
$$\partial_t n_n + \nabla \cdot (n_n V_{0,n}) = -n_e n_n K_{ioniz}$$

A 2D z - θ test case to observe the ECDI ¹⁾

- ◆ Extracted from a z - θ collisionless PIC simulation with given ionization rate ¹⁾



Given magnetic field and ionization source term for the test case



$$J = e \int_0^d S(x) dx = 400 \text{ A/m}^2$$

- ◆ Applied voltage of 200V
- $n_{e,init} = n_{i,init} = 5 \cdot 10^{16} \text{ m}^{-3}$
- ◆ $T_{e,init} = 5 \text{ eV} = 58000 \text{ K}$; $T_{i,init} = 300 \text{ K}$

Goal : Study this test case with our fluid formulation and observe if the ECDI is triggered by fluid equations

A 2D z - θ test case to observe the ECDCI ¹⁾

A simplified fluid model :

$$\begin{aligned}\partial_t n_e + \nabla \cdot (n_e u_e) &= S, & \partial_t n_i + \nabla \cdot (n_i u_i) &= S \\ \partial_t (m_e n_e V_e) + \nabla \cdot (m_e n_e V_e V_e + k_B T_e n_e I) &= -en_e (E + V_e \times B) \\ \partial_t (m_i n_i V_i) + \nabla \cdot (m_i n_i V_i V_i + k_B T_i n_i I) &= en_i (E + V_i \times B) \\ \partial_t (m_e n_e E_e) + \nabla \cdot \left(\left(\frac{1}{2} m_e n_e V_e^2 + \frac{\gamma}{\gamma - 1} P_e \right) \cdot V_e \right) &= -en_e E \cdot V_e + \Theta_e \\ \partial_t (m_i n_i E_i) + \nabla \cdot \left(\left(\frac{1}{2} m_i n_i V_i^2 + \frac{\gamma}{\gamma - 1} P_i \right) \cdot V_i \right) &= en_i E \cdot V_i + \Theta_i\end{aligned}$$

How to express energy source terms ?

- An ion is created at $T_i = 0.5$ eV

$$\Theta_i = S k_B T_i$$

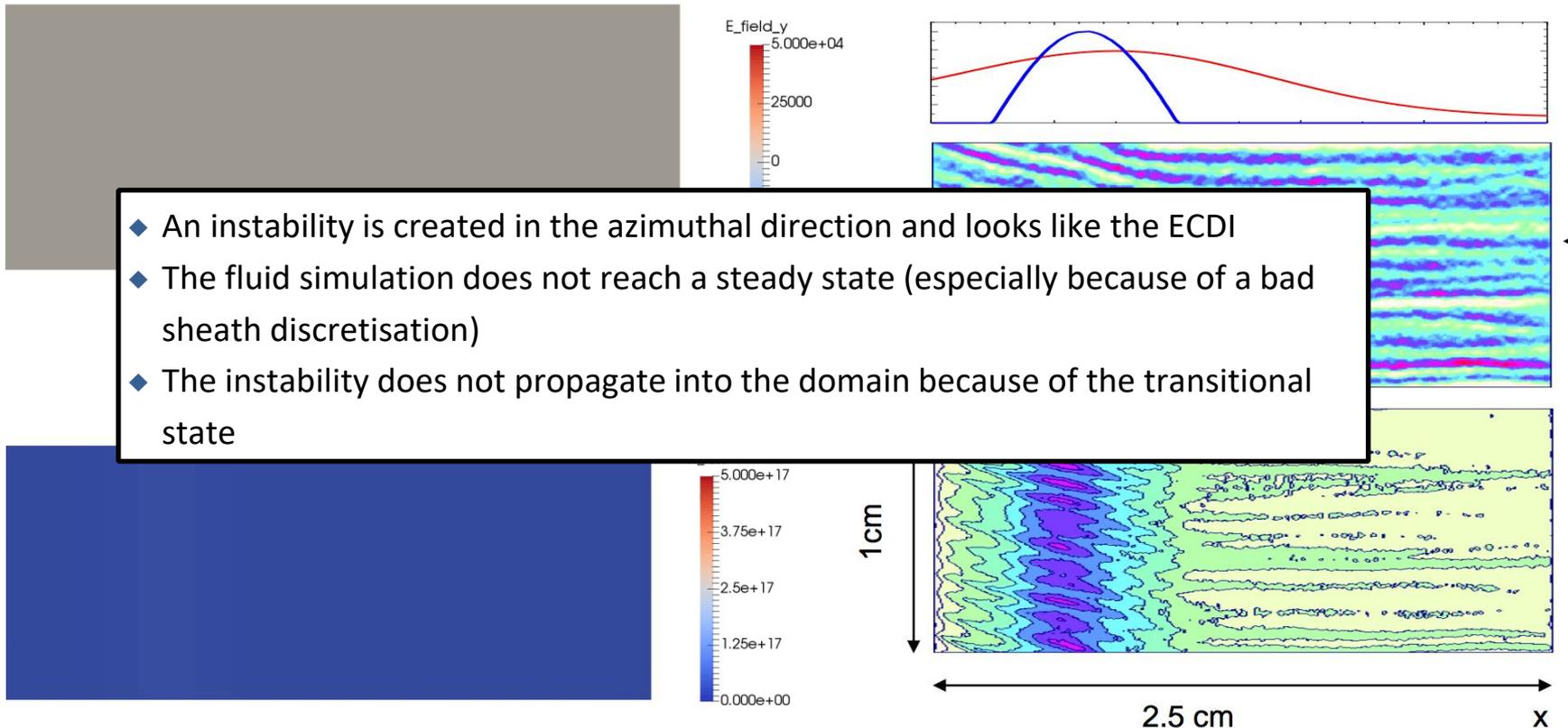
- An electron is created at $T_e = 10$ eV

$$\Theta_e = S k_B T_e$$

A 2D z - θ test case to observe the ECDI ¹⁾

Simulation 1 with a 500x200 mesh (1 cell per λ_D) and a second order scheme (HLLC-MUSCL) :

Mean azimuthal electric field

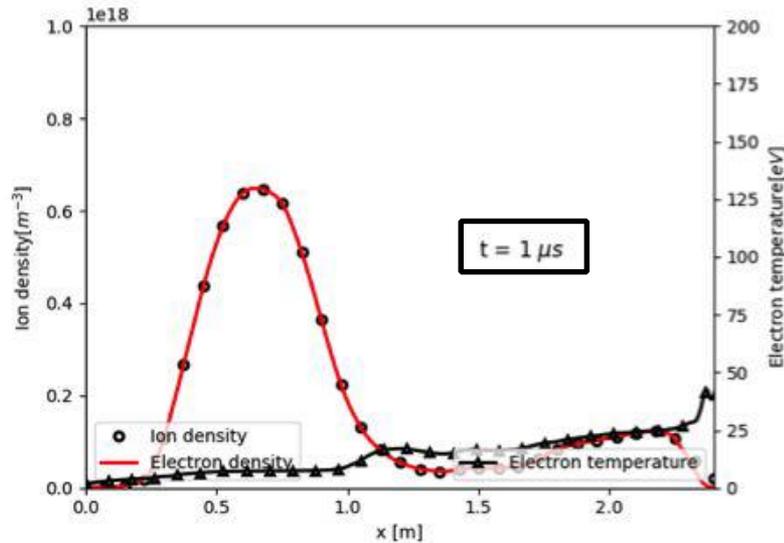


- ◆ An instability is created in the azimuthal direction and looks like the ECDI
- ◆ The fluid simulation does not reach a steady state (especially because of a bad sheath discretisation)
- ◆ The instability does not propagate into the domain because of the transitional state

Time: $0.01 \mu\text{s}$

Azimuthal electric field and ion density at steady state⁽¹⁾

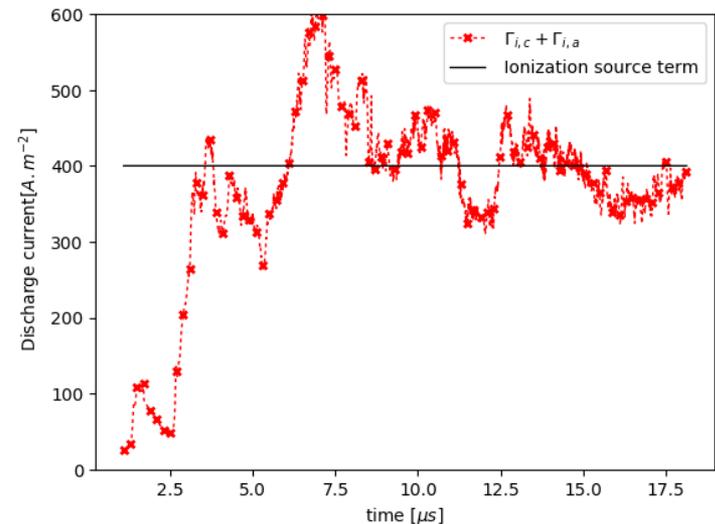
A 2D z- θ test case to observe the ECDI ¹⁾



Azimuthal mean electron and ion density and electron temperature during the transitional state

- ◆ The plasma density is too high in the ionization region
- ◆ The bad discretisation of the sheath creates a high electron temperature which pollutes the entire domain
- ◆ The sum of ion currents oscillates around the value of ionization current

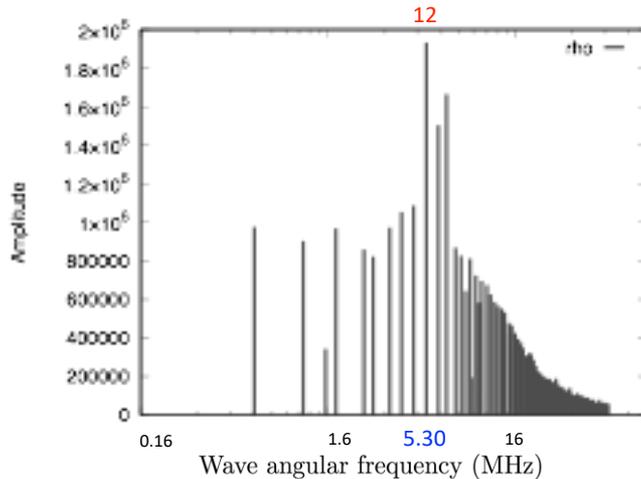
$$\Gamma_{i,c} + \Gamma_{i,a} = e \int_0^d S(x) dx = 400 \text{ A/m}^2$$



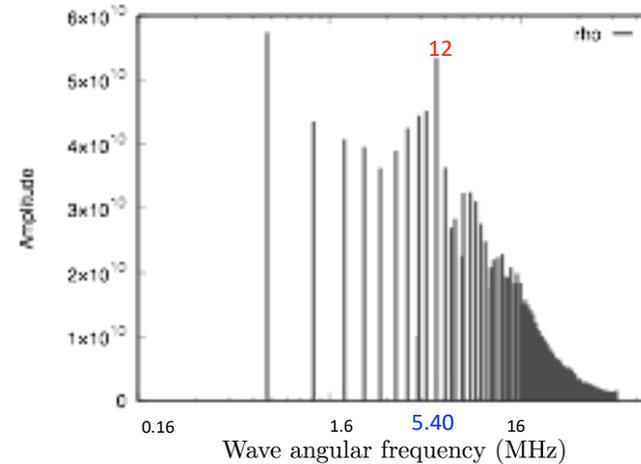
Comparison of the sum of the ion current through electrodes vs creation of ions by ionization

A 2D z - θ test case to observe the ECDI ¹⁾

Simulation 1 with a 500x200 mesh (1 cell per λ_D) and a second order scheme (HLLC-MUSCL) :



Ion density spectrum of the simulation 1



Azimuthal electric field spectrum of the simulation 1

Theoretical results for the ECDI :

$$\lambda_\omega = 2\pi\sqrt{2}\lambda_{De} \approx 1\text{mm}$$
$$v_\omega = c_s\sqrt{\frac{2}{3}} \approx 5\text{km/s}$$
$$f_\omega = \frac{v_\omega}{\lambda_\omega} \approx 5\text{MHz}$$

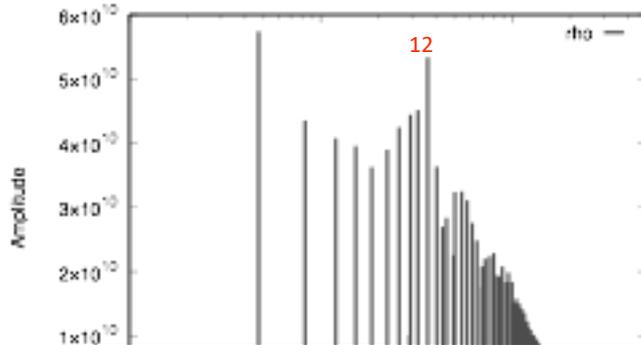
- ◆ The 12th mode seems to correspond to the good frequency and is one of the highest mode



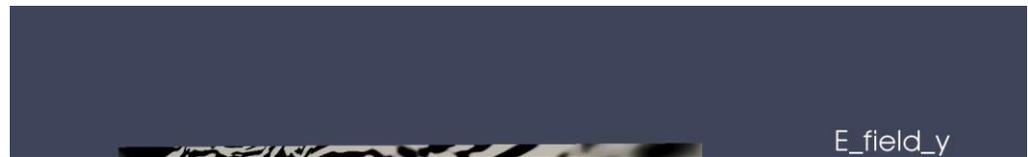
Investigate this particular mode using a Domain Decomposition Method

A 2D z - θ test case to observe the ECDI ¹⁾

Simulation 1 with a 500x200 mesh (1 cell per λ_D) and a second order scheme (HLLC-MUSCL) :



Electric field 12th mode from the Domain Mode Decomposition



- ◆ The instability of the 12th mode is slower with a smaller wavelength but at the same frequency as the ECDI
- ◆ The amplitude of the instability is high in the region of strong magnetic field
- ◆ However the ECDI is perturbed by the transitional regime due to the initialisation

$$\lambda_\omega = 2\pi\sqrt{2}\lambda_{De} \approx 1mm$$

$$v_\omega = c_s\sqrt{\frac{2}{3}} \approx 5km/s$$

$$f_\omega = \frac{v_\omega}{\lambda_\omega} \approx 5MHz$$



Time: 2.00 μ s

$$\lambda_{mode_{12}} \approx 0.5mm$$

$$v_{mode_{12}} \approx 2.7km/s$$

Conclusions

- An instability appears in the azimuthal direction which has same properties as the Electron Cyclotron Drift Instability
- The fluid simulation fails to reach a steady state, possibly due to the cathode injection model

- Note that a second order numerical scheme is mandatory for a good resolution of electrode sheaths¹⁾